

X-RAY AND OPTICAL SPECTRAL PROPERTIES OF
ACTIVE GALACTIC NUCLEI

by

LAURA ELIZABETH TROUILLE

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*Il est encore plus facile de juger
de l'esprit d'un homme
par ses questions que par ses réponses.*

- Voltaire

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Contents

Acknowledgements	ii
List of Tables	vii
List of Figures	ix
1 Introduction	1
1.1 Motivation	2
1.2 X-ray versus Optical Spectral Properties of AGNs	2
1.3 Is [OIII] a Reliable Indicator of AGN Activity?	6
1.4 The Elusive Population of Compton-thick AGNs	7
1.5 The OPTX Survey Fields	10
1.5.1 CDF-N	10
1.5.2 CLASXS & CLANS	10
1.6 Thesis Goals and Outline	10
References	11
2 The OPTX Project I: The Flux and Redshift Catalogs for the CLANS, CLASXS, and CDF-N Fields	16
Abstract	17
2.1 Introduction	18
2.2 X-ray Properties	20
2.2.1 CLANS X-ray Observations	20
2.2.2 CLANS X-ray Source Detection and Fluxes	21
2.2.3 CLANS X-ray Catalog	23
2.2.4 Flux Limits	23

2.2.5	Effective Areas	24
2.3	Optical and IR Imaging	27
2.3.1	MegaCam	27
2.3.2	WIRCam	28
2.3.3	ULBCam	29
2.3.4	WFCAM	29
2.3.5	Spitzer	30
2.3.6	Optical and NIR Source Detection and Photometry	30
2.3.7	CLASXS: Updated Optical Photometry	32
2.4	Redshift Information	32
2.4.1	Spectroscopic Observations	32
2.4.2	Spectroscopic Completeness	34
2.4.3	Photometric Redshifts	34
2.4.4	Redshift and Flux Distributions	37
2.5	Optical Spectral Classification	42
2.6	Optical/IR Counterparts Catalogs	44
2.6.1	CLANS	44
2.6.2	CLASXS	44
2.6.3	CDF-N	45
2.7	X-ray Luminosities	45
2.8	Number Counts	47
2.9	Summary	52
	References	52
3	The OPTX Project III: X-ray Versus Optical Spectral Type for AGNs	56
	Abstract	57
3.1	Introduction	58
3.2	X-ray Data	60
3.3	Classification by Optical Spectral Type	61
3.4	Comparison of Optical Classification with X-ray Classification	64

3.4.1	Γ_{eff} Distributions	64
3.4.2	Γ_{eff} versus X-ray Luminosity	69
3.5	Discussion	72
3.6	Conclusions	74
	References	74
4	The OPTX Project IV: How Reliable is [OIII]λ5007 as a Measure of AGN Activity?	78
	Abstract	79
4.1	Introduction	80
4.2	Sample	81
4.3	High-Ionization Narrow Emission Lines	82
4.3.1	[OIII] λ 5007	82
4.3.2	Emission-Line Reddening	84
4.4	Emission-Line Ratio Diagnostic Diagrams	86
4.5	The Relation Between $L_{[\text{OIII}]}$ and L_X	87
4.6	Discussion	92
4.7	Summary	96
	References	98
5	The OPTX Project V: Identifying Distant Compton-thick AGNs	101
	Abstract	102
5.1	Introduction	103
5.2	Sample	104
5.2.1	SDSS	104
5.2.2	OPTX	105
5.2.3	GOODS-N	106
5.2.4	Rest-frame $g - z$ Color	106
5.3	High-Redshift Optical Emission Line Diagnostic	106
5.4	Identifying X-ray Selected AGNs Using Optical Emission Line Diagnostics	109
5.5	Identifying X-ray Faint/Undetected AGNs using Optical Emission Line Diagnostics	109

5.6	Application: Identifying Compton-thick AGNs	110
5.7	Discussion	114
5.7.1	A New Diagnostic Diagram for Galaxies at Redshifts up to 1.4	114
5.7.2	CT AGN Space Densities compared with X-ray Background Models	116
5.8	Summary	117
	References	119
6	Conclusions	122
6.1	Thesis Results	123
6.2	Completion of OPTX Project V	126
	References	129

List of Tables

2.1	Spectroscopic Completeness of Selected Surveys from the Literature	19
2.2	CLANS Observation Summary	20
2.3	X-ray Data Specifications	21
2.4	CLANS X-ray Catalog: Basic Source Properties	23
2.5	CLANS X-ray Catalog: Additional Source Properties	24
2.6	CLANS Optical & Near-IR Specifications	31
2.7	CLASXS Optical & Near-IR Specifications	32
2.8	CDF-N Near-IR Specifications	33
2.9	CLASXS Zeropoint Offsets	34
2.10	X-ray Samples by Field and Optical Spectral Class	42
2.11	CLANS Optical/IR Counterparts Catalog	43
2.12	CLASXS Optical/IR Counterparts Catalog	43
2.13	CDF-N Optical/IR Counterparts Catalog	44
2.14	Best-fit to the Number Counts	47
3.1	CLANS, CLASXS, and CDF-N Survey Characteristics	61
3.2	CLANS Catalog, Updated Sources	62
3.3	Optical Spectral Classification for our 2 – 8 keV Sample	63
3.4	Non-BLAGNs with Soft Photon Indices	71
4.1	$\log(L_{[\text{OIII}]} / L_X)$	89
4.2	Best fits to the luminosity dependent ratio of $L_{[\text{OIII}]} / L_X$	91
5.1	GOODS-N $z < 0.5$ Candidate AGNs	110

5.2	GOODS-N $z < 1.4$ Candidate AGNs	113
5.3	CT AGN Space Density Determinations from the Literature	118

List of Figures

1.1	Composite Spectra of Broad-Line and non-Broad-Line AGNs	3
1.2	Simple Unified Model for AGNs	4
1.3	Redshift Bias in X-ray Classification	5
1.4	Clumpy Torus and Disk-Wind Models for AGN Obscuration	6
1.5	Observed Extragalactic X-ray Background	8
2.1	Location of the CLANS and CLASXS pointings	22
2.2	Flux vs. Off-axis Angle for the OPTX Sources	25
2.3	Probability of Source Detection	27
2.4	Fraction of Annulus within ACIS-I Chip	28
2.5	Total Effective Area	29
2.6	Spectroscopic Completeness	35
2.7	Template SEDs	37
2.8	Reliability of Photo-z Determinations	38
2.9	Comparison of Template Fitting and HYPERZ Methods	39
2.10	OPTX Redshift Distribution	39
2.11	OPTX 0.5 – 2 keV Flux Distribution	40
2.12	OPTX 2 – 8 keV Flux Distribution	41
2.13	X-ray Luminosity vs. Redshift	46
2.14	Differential Number Counts	49
2.15	Differential Number Counts, Residuals	50
2.16	Comparison of Number Counts Fits from Other Surveys	51

3.1	2 – 8 keV Flux Distribution	62
3.2	R Magnitude vs. 2 – 8 keV Flux	64
3.3	Spectroscopic Completeness	65
3.4	Redshift Distribution by Optical Spectral Type	66
3.5	Γ_{eff} Distribution by Optical Spectral Type	67
3.6	$I - J$ vs. $J - K$ for our X-ray Hard BLAGNs	68
3.7	Γ_{eff} vs. X-ray Luminosity	70
3.8	Γ_{eff} vs. X-ray Luminosity in Redshift Bins	72
4.1	L_X vs. Redshift for our 2 – 8 keV Sample	83
4.2	Emission-line Reddening Correction	84
4.3	BPT Diagram Applied to our 2 – 8 keV Sample	87
4.4	Distribution of $L_{[\text{OIII}]} / L_X$ for our 2 – 8 keV Sample	88
4.5	$L_{[\text{OIII}]} / L_X$ vs. L_X for our OPTX Non-BLAGNs	90
4.6	$L_{[\text{OIII}]} / L_X$ vs. L_X for our OPTX BLAGNs	91
4.7	BPT Diagram, Increasing $L_{[\text{OIII}]} / L_X$	93
4.8	$L_{[\text{OIII}]} / L_X$, Sources with Reliable Balmer Fluxes	94
4.9	$L_{[\text{OIII}]} / L_X$ vs. L_X , BLAGN Comparison	95
4.10	Comparison of X-ray and OIII LFs	97
5.1	BPT Diagram Applied to the SDSS DR7 Galaxies	107
5.2	High-redshift Diagnostic Applied to the SDSS DR7 Galaxies	108
5.3	High-redshift Diagnostic Applied to our OPTX Non-BLAGNs	109
5.4	BPT Diagram Applied to the GOODS-N Galaxies	111
5.5	High-redshift Diagnostic Applied to the GOODS-N Galaxies	112
5.6	Distribution in $L_{[\text{NeIII}]} / L_X$ for our OPTX Non-BLAGNs	113
5.7	Location of LINERs in the BPT Diagram	114
5.8	Comparison of CT AGN Space Densities from the Literature	117
6.1	BPT Diagram for the CLANS Galaxy Survey	128

Chapter 1

Introduction

1.1 Motivation

Active Galactic Nuclei (AGNs) are cosmological powerhouses that emit prodigious amounts of light over several decades in energy. They number among the most luminous radio, optical, and X-ray sources known. This relatively uniform, luminous, multi-band emission from AGNs is non-stellar in origin. AGN unification models suggest that all AGNs are powered by accretion onto a supermassive black hole (SMBH; $M_{BH} \geq 10^6 M_{\odot}$) that resides at the center of the galaxy (see Antonucci 1993 for a review).

Current observations suggest that most galaxies, including the Milky Way, have a SMBH at their centers (Kormendy & Richstone 1995; Magorrian et al. 1998; van der Marel 1999). Moreover, the mass of the SMBH appears to be tied to the mass of the host galaxy’s bulge from studies of both velocity dispersion (Ferrarese & Merritt 2000; Gebhardt et al. 2000) and the bulge luminosity (Magorrian et al. 1998). This strongly suggests that the formation of the host galaxy and its SMBH are somehow connected.

Hard 2 – 8 keV X-ray surveys provide as complete and unbiased a sample of AGNs as is presently possible (Mushotzky et al. 2004). At these energies, dilution of the AGN light by the stellar light from the host galaxy is minimal, and the cross-section of the circumnuclear obscuring material drops quickly with increasing X-ray energy, enabling hard X-ray photons to penetrate all but the densest material. X-ray surveys have so far found the highest density of AGNs on the sky ($\geq 5500 \text{ deg}^{-2}$, as compared to $\sim 400 \text{ deg}^{-2}$ for the deepest optical surveys). Ultradeep *Chandra* X-ray surveys (Mushotzky et al. 2000; Brandt et al. 2001; Giacconi et al. 2002; Alexander et al. 2003; Barger et al. 2005; Luo et al. 2008) have resolved nearly 100% of the 2 – 8 keV X-ray background (XRB; see Churazov et al. 2007 for a recent measurement of the XRB using INTEGRAL and Gilli et al. 2007 and Frontera et al. 2007 for in-depth comparisons of the various XRB measurements to date). Even though the majority of the 0.5 – 8 keV XRB has been resolved into discrete sources, the nature of the sources is still unfolding.

The increased emergence of wide-field, multi-object spectrographs like HYDRA (Barden et al. 1994) on the WIYN 3.5 m telescope and the DEep Imaging Multi-Object Spectrograph (DEIMOS; Faber et al. (2003) on the Keck 10 m telescope on Mauna Kea greatly improves the efficiency of spectroscopic surveys designed to examine the optical properties and measure redshifts for these X-ray detected AGNs. The combination of X-ray and optical data allows one to study the X-ray/optical AGN connection and its evolution over cosmic time.

1.2 X-ray versus Optical Spectral Properties of AGNs

There are striking differences in the optical spectra of AGNs. Some show broad (FWHM $> 2000 \text{ km s}^{-1}$) emission lines while others exhibit narrow (FWHM $\sim 400 \text{ km s}^{-1}$) forbidden emission lines (see Figure 1.1). One generally assumes that there is less intrinsic diversity among AGNs than we observe, and that the variety of AGN phenomena is due to a combination of real differences in a small number of physical parameters (like luminosity) coupled with apparent differences which are due to observer-dependent parameters (like orientation). In the Antonucci (1993) simple unified model for AGNs, it is the geometry of a dust torus which determines the amount of obscuring material along the observer’s line of sight to the central X-ray emitting regions, as well as which region (broad-line or narrow-line) one is observing. In this

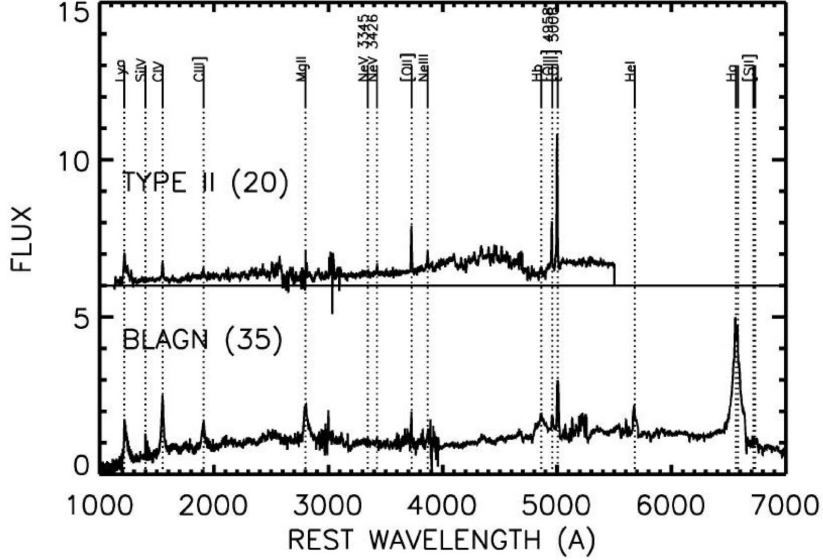


Fig. 1.1. — Composite spectra of 35 broad-line AGNs and 20 non-broad-line AGNs, reproduced from Barger et al. (2005).

model, broad-line AGNs are unobscured sources in which one is seeing down the ionization cone to clouds in high density regions within ~ 0.003 parsecs of the accretion disk. Non-broad-line AGNs, on the other hand, are observed when the central source is obscured by the circumnuclear material and the emission lines come from lower density clouds at 30 – 100 parsecs from the accretion disk (see Figure 1.2 showing the Urry & Padovani 1995 cartoon depiction of the simple unified model for AGNs).

X-ray data alone have also long been used to estimate the amount of obscuration between the observer and the nuclear source through the 0.5 – 8 keV spectral slope. Since 2 – 8 keV X-rays will penetrate obscuring material (except in Compton-thick AGNs, where the neutral hydrogen column density, N_H , in the line of sight is higher than the inverse Thomson cross section, $N_H = 1.5 \times 10^{24} \text{ cm}^{-2}$) and 0.5 – 2 keV X-rays will not, in low signal-to-noise data a shallower slope may indicate an obscured source. X-ray spectra can be approximated with a power-law of the form $P(E) = AE^{-\Gamma_{\text{eff}}}$, where E is the photon energy in keV and A is the normalization factor. (Here, we indicate the power-law slope as Γ_{eff} to distinguish it from the intrinsic slope, Γ . Γ_{eff} is not the true Γ unless there is no intrinsic absorption.)

Hasinger et al. (2005) initiated a current trend in the field to combine the X-ray and optical classification schemes in order to create the most ‘complete’ sample of unobscured AGNs. Unobscured AGNs, in this scenario, include any object optically classified as a broad-line AGN, as well as any object satisfying $L_X \geq 10^{42} \text{ erg s}^{-1}$ and $\Gamma_{\text{eff}} \geq 1.2$. This decision was based on their assumption that ‘true’ broad-line

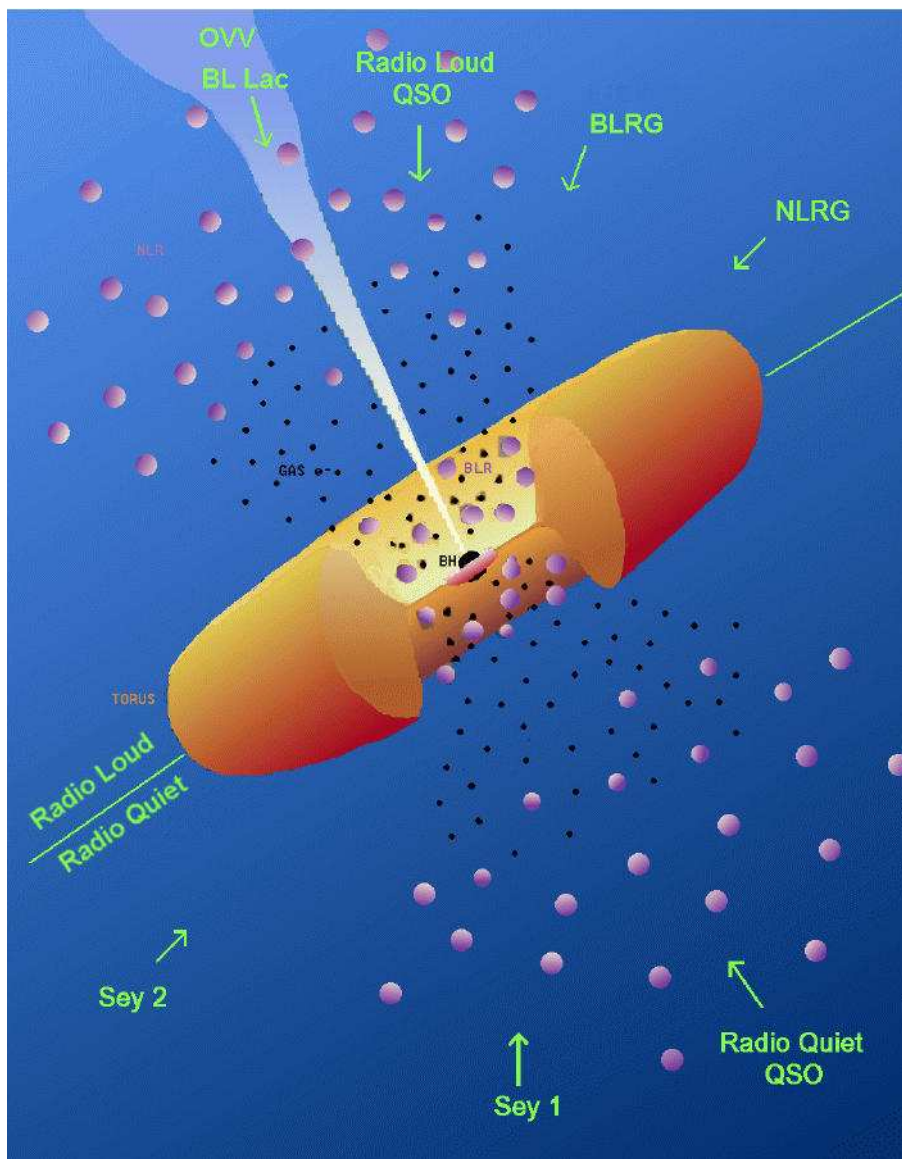


Fig. 1.2. — Cartoon depiction of the simple unified model for AGNs, reproduced from Urry & Padovani (1995).

AGNs may be optically misclassified as non-broad-line AGNs due to dilution of the AGN light by the host galaxy light, particularly at low luminosities (Moran et al. 2002; Severgnini et al. 2003; Garcet et al. 2007; Cardamone et al. 2007).

The creation of a classification scheme that mixes optical and X-ray spectral diagnostics requires a thorough understanding of the correspondence between X-ray and optical spectral type. Unfortunately, 10 – 30% of AGNs have (1) X-ray spectra that show no absorption and (2) optical spectra that suggest obscuration (e.g., Pappa et al. 2001; Panessa & Bassani 2002; Barcons et al. 2003; Georgantopoulos & Zezas 2003; Caccianaga et al. 2004; Corral et al. 2005; Wolter et al. 2005; Tozzi et al. 2006). The opposite effect, i.e., (1) X-ray spectra that show absorption and (2) optical spectra that suggest no obscuration, has also

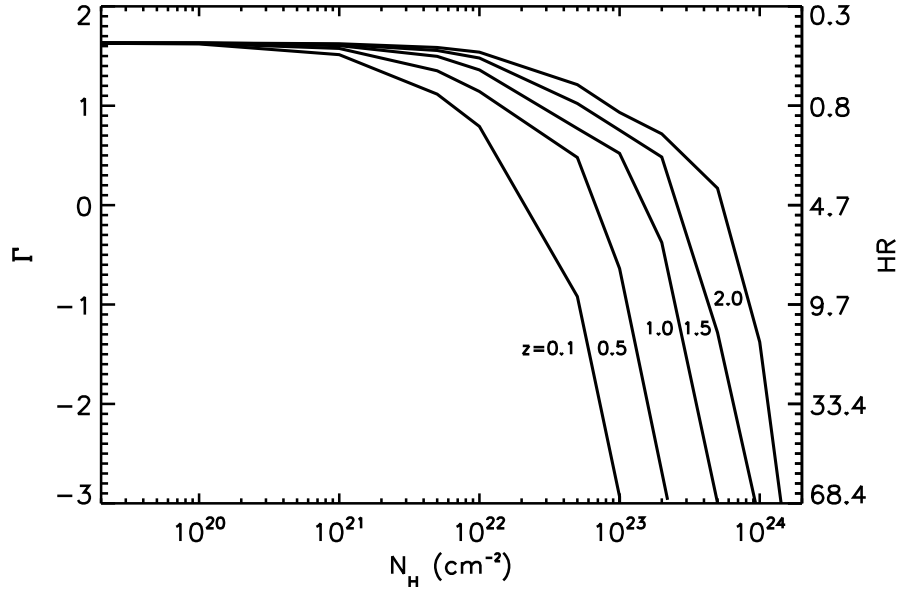


Fig. 1.3. — Γ_{eff} vs. N_H for $z = 0.1, 0.5, 1.0, 1.5,$ and 2.0 , as labeled. The right-hand y-axis provides the corresponding values for the *Chandra*-specific hardness ratio (2–8 keV count rate divided by the 0.5–2 keV count rate). Values determined using WEBPIMMS, assuming a spectral index of 1.7 and the Galactic HI column density in the direction of the Lockman Hole ($5.7 \times 10^{19} \text{ cm}^{-2}$).

been observed (e.g., Comastri et al. 2001; Wilkes et al. 2002; Fiore et al. 2003; Brusa et al. 2003; Akiyama et al. 2003; Silverman et al. 2005; Gallagher et al. 2006; Hall et al. 2006; Tajer et al. 2007). Furthermore, all X-ray based classification schemes suffer from a serious redshift bias (as one observes higher redshift sources, the 2–10 keV filter samples higher energies which are less affected by obscuring material, and the source appears softer; Szokoly et al. 2004). Figure 1.3, showing Γ_{eff} versus N_H for a range of redshifts, illustrates this effect – a higher redshift source with the same amount of obscuration has a higher value of Γ_{eff} .

The sources which show a discrepancy in their X-ray and optical indicators of absorption along the line-of-sight are not only problematic for mixed classification schemes, but also in terms of the validity of the simple unified model for AGNs (i.e., the homogeneous dust ‘doughnut’ surrounding the central engine). An important question is how many of the X-ray unabsorbed, optically obscured sources can be explained by the often invoked observational complications like host galaxy dilution (Moran et al. 2002; Severgnini et al. 2003; Hasinger et al. 2005), X-ray variability (Paolillo et al. 2004), and the redshifting of broad lines out of the spectral window (Silverman et al. 2005). If such sources cannot be explained as suffering from these observational complications, they lend support for the clumpy torus model (Risaliti et al. 2000; Hönig et al. 2006; Elitzur 2007; Nenkova et al. 2008) and torus-as-a-wind scenarios (Emmering et al. 1992; Konigl & Kartje 1994; Bottorff et al. 1997; Kartje et al. 1999). In the clumpy torus models, the probability for direct viewing of the central engine decreases away from the axis but is always finite (see Figure 1.4). These models also incorporate the change in cloud composition across the dust sublimation radius. While dusty material absorbs continuum radiation both in the UV/optical and X-ray, dust-free gas attenuates just the X-ray continuum. Clouds inside the dust sublimation radius therefore only provide obscuration in the X-ray regime. The combination of clumpiness and radius dependent dust-to-gas ratios allows for sources with differing X-ray and

optical spectral type.

A highly optically spectroscopically complete sample of X-ray selected AGNs will allow us to test the validity of the mixed classification scheme and reliably determine the percentage of AGNs whose X-ray properties indicate significant absorption but whose optical spectra show little obscuration, and vice-versa. We will be able to quantify how many sources with discrepant X-ray and optical properties can be explained by observational complications (like host galaxy dilution, X-ray variability, and the redshifting of broad lines), and whether a significant percentage truly do call for alterations to the simple unified model for AGNs.

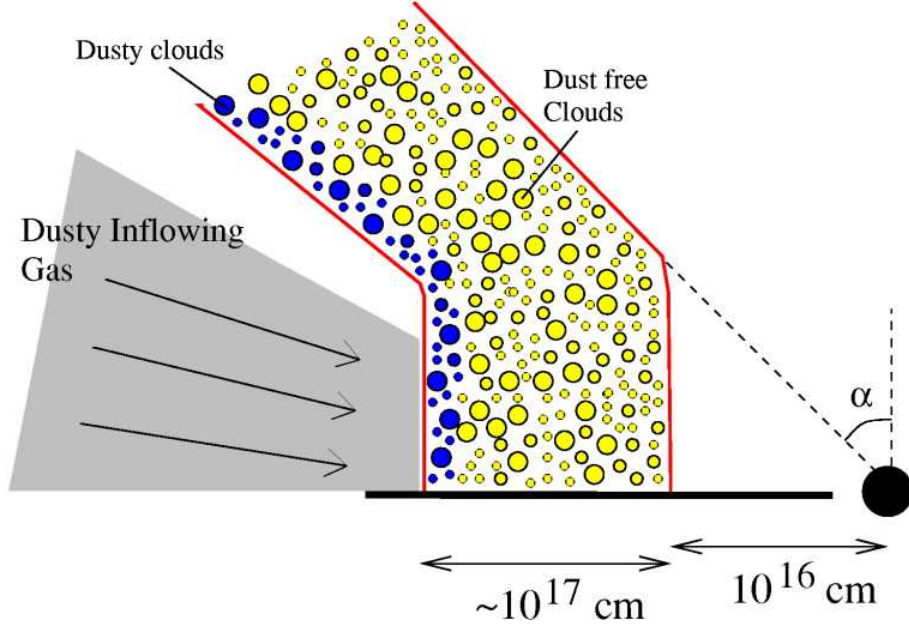


Fig. 1.4. — Clumpy torus model, derived from Elvis (2000) and further developed in Nenkova et al. (2008, and references therein). Reproduced from Risaliti et al. (2002).

1.3 Is [OIII] a Reliable Indicator of AGN Activity?

Kauffmann et al. (2003) and Heckman et al. (2004) proposed using the luminosity of the [OIII] λ 5007Å (hereafter, [OIII]) line as a tracer of AGN activity, thus efficiently selecting AGNs from the large sample of galaxies in the Sloan Digital Sky Survey (SDSS). Narrow emission lines are thought to arise in gas excited by ionizing radiation escaping along the polar axis of the obscuring ‘torus’. Since the narrow line region lies outside of the dusty circumnuclear material, the lines should not suffer from obscuration by that high column density material (though they may be affected by dust within the host galaxy). The [OIII] line is particularly appealing because it is one of the strongest narrow emission lines in optically obscured AGNs and the least contaminated by contributions from HII regions associated with star formation in the host galaxy (Heckman et al. 2005). While there is some evidence that the [OIII] luminosity is a good measure of AGN activity (e.g. Mulchaey et al. 1994), and, indeed, many researchers have assumed it to be so in their analyses (e.g. Alonso-Herrero et al. 1997; Hao et al. 2005; Netzer et al. 2006; Bongiorno et al. 2010), recent work has called this conclusion into question (e.g., Cocchia et al. 2007; Meléndez et al. 2008; Diamond-Stanic

et al. 2009; La Massa et al. 2009).

AGN activity can also be traced through the 2 – 10 keV luminosity, although this measurement may suffer from some absorption (e.g., Diamond-Stanic et al. 2009; La Massa et al. 2009). Heckman et al. (2005) used their measurements of the ratio between X-ray and [OIII] luminosity to transform the Sazonov et al. (2004) *Rossi X-ray Timing Explorer (RXTE)* X-ray luminosity function (LF) into an [OIII] LF. They then compared this [OIII] LF with the SDSS [OIII] LF from Hao et al. (2005). They found that at the same equivalent luminosity, the space density of the hard X-ray selected AGNs was lower than that of the [OIII]-selected AGNs by a factor of ~ 0.5 dex. This led them to argue that selection by hard X-rays misses a significant fraction of the local AGN population with strong emission lines.

A large sample of X-ray selected AGNs with high optical spectroscopic completeness is ideal for investigating the reliability of the [OIII] emission line as an indicator of AGN activity. Such a sample can also test how many X-ray selected AGNs are misidentified as star formers by the commonly used Baldwin et al. (1981)-type narrow emission-line ratio empirical diagnostic diagram ([OIII]/H β versus [NII]/H α , hereafter BPT diagram; see also Osterbrock et al. 1985 and Veilleux et al. 1987). Finally, redoing the Heckman et al. (2005) X-ray and [OIII] LF comparison with (1) a *Chandra* X-ray selected sample of AGNs combined with the *SWIFT BAT* results and (2) an updated understanding of the relation between the [OIII] and X-ray luminosities, will allow a more accurate estimate of the efficiency of optical versus X-ray selection of AGNs.

1.4 The Elusive Population of Compton-thick AGNs

The extragalactic XRB is thought to be mainly composed of the emission from accreting SMBHs throughout the history of the universe. The energy density of this cosmic radiation resides predominantly in the hard X-ray spectrum, which peaks at 30 keV. From the number density of known AGNs, we can estimate their contribution to the XRB. Optically bright quasars and Seyfert galaxies dominate at low energies (up to a few keV), while obscured AGNs, which outnumber unobscured ones by a factor of 3 – 4 (Ueda et al. 2003; La Franca et al. 2005), are responsible for the bulk of the XRB at high energies (> 10 keV). The ultradeep *Chandra* X-ray surveys have resolved nearly 100% of the 2 – 8 keV XRB, but we appear to be missing up to 30% of the XRB at 30 keV (Gilli et al. 2007). The missing intensity of the XRB around 30 keV is attributed to the integrated emission from the elusive population of heavily X-ray absorbed, Compton-thick (CT) AGNs (with $N_H > 1.5 \times 10^{24}$ cm $^{-2}$). This missing population represents a significant fraction of the cosmic accretion onto SMBHs (Marconi et al. 2004).

This estimate for the number density of CT AGNs depends on (1) the true absolute intensity of the hard XRB, which is uncertain at the 10 – 30% level, (2) the average broadband spectra of Compton-thin AGNs, and (3) the cosmological evolution of CT AGNs, which is assumed to be the same as Compton-thin AGNs.

After the first pioneering XRB measurements (Horstman et al. 1975), the *High Energy Astronomical Observatory (HEAO-1)* carried out a major effort in the late 1970s to obtain a reliable estimate of the XRB spectrum from 2 – 400 keV (Marshall et al. 1980; Gruber et al. 1999). More recent experiments with the focusing telescopes aboard *BeppoSAX* (Vecchi et al. 1999), *XMM-Newton* (Lumb et al. 2002; De Luca & Molendi 2004), and *Chandra* (Hickox & Markevitch 2006) determined the 2 – 10 keV XRB to be $\sim 40\%$

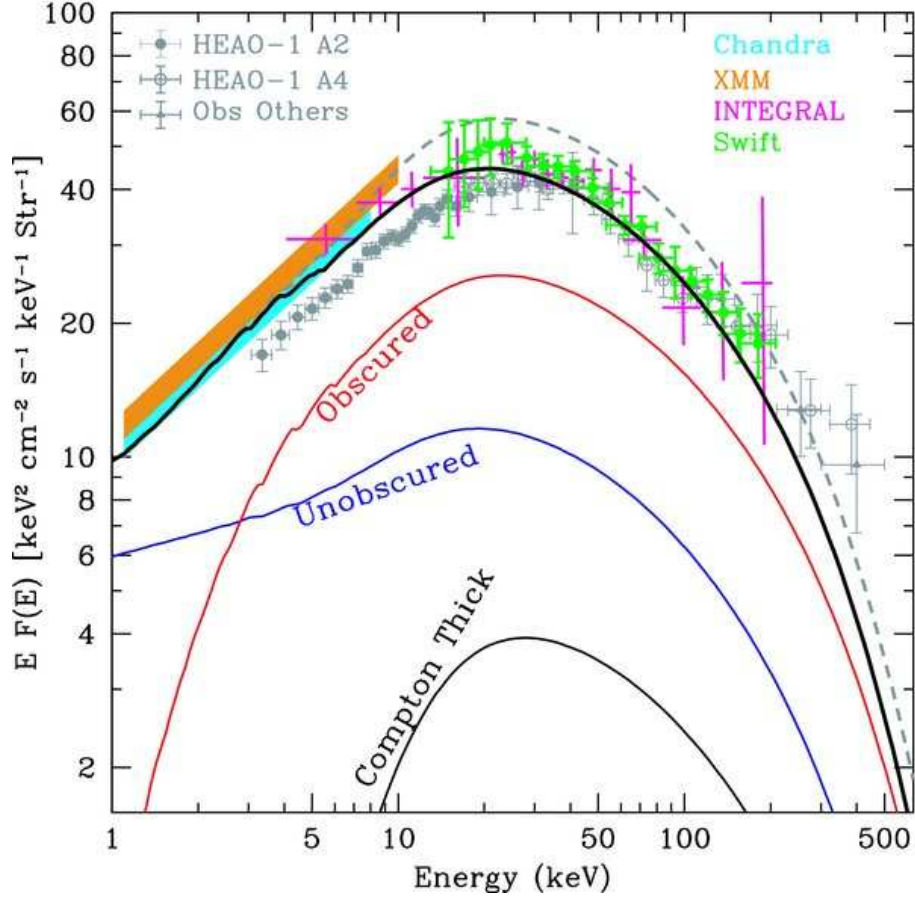


Fig. 1.5. — Observed spectrum of the extragalactic XRB from *HEAO-1* (Gruber et al. 1999), *Chandra* (Hickox & Markevitch 2006), *XMM-Newton* (De Luca & Molendi 2004), *INTEGRAL* (Churazov et al. 2007), and *Swift* (Ajello et al. 2008) data. The dashed gray line shows the XRB spectrum from the AGN population synthesis model of Treister et al. (2005), which assumed a 40% higher value for the *HEAO-1* XRB normalization. The thick black solid line shows the Treister et al. (2009) new population synthesis model for the XRB spectrum; the only change is the number of CT AGNs, which is reduced by a factor of 4 relative to the number in Treister et al. (2005). Reproduced from Treister et al. (2009).

higher than the *HEAO-1* estimate. Driven by these discordant results, several authors (Ueda et al. 2003; De Luca & Molendi 2004; Comastri 2004; Worsley et al. 2005; Worsley et al. 2006; Ballantyne et al. 2006; Hopkins et al. 2006), in their evaluation of the fraction of the XRB that can be resolved into individual sources, increased the *HEAO-1* XRB intensity by a factor of ~ 1.3 over the entire energy band (driving up the required number density of CT AGNs). Such a high value of the XRB intensity at 10 – 50 keV is now ruled out by new *INTEGRAL* (Churazov et al. 2007), *Swift* (Ajello et al. 2008), and *BeppoSAX* (Frontera et al. 2007) data.

Treister et al. (2009) used these updated XRB measurements to create a new AGN population synthesis model with a more realistic prediction for the number density of CT AGNs. Figure 1.5 shows the results of this new population synthesis model (thick black line) compared to the measurements of the XRB by *HEAO-1*, *Chandra*, *XMM*, *INTEGRAL*, and *Swift*. The thin black line shows the predicted contribution to

the XRB from CT AGNs.

Assuming that the most recent XRB measurements are correct and that the new AGN population synthesis models are valid, one can then determine if there remains a significant population of missing sources by comparing the predicted number density of CT AGNs with the observed number density of CT AGNs.

The *Swift* and *INTEGRAL* ultra-hard X-ray detectors (> 15 keV) are sensitive to the direct emission from the central engine in transmission-dominated CT AGNs ($N_H \approx 10^{24} - 10^{25}$ cm $^{-2}$) and have helped constrain the observed number density of CT AGNs at very bright fluxes ($f_{17-70\text{keV}} > 10^{-11}$ erg cm $^{-2}$ s $^{-1}$) (Beckmann et al. 2006; Bassani et al. 2006; Sazonov et al. 2007; Winter et al. 2008; Winter et al. 2009; Tueller et al. 2010). These surveys, however, currently only probe out to $z < 0.05$ and miss the higher column density ($N_H > 10^{25}$ cm $^{-2}$), reflection-dominated CT AGNs.

Hard X-ray selection by *Chandra* or *XMM-Newton*, while being the most direct means currently to find obscured sources at high redshift, is limited by telescope sensitivities and wavelength coverage. Only the brightest sources with non-extreme column densities can be identified (Brandt & Hasinger 2005). For example, the archetype CT AGN NGC 1068, with $L_X = 10^{44}$ erg s $^{-1}$, would have a hard X-ray flux of only 7×10^{-18} erg cm $^{-2}$ s $^{-1}$ at $z = 2$, more than an order of magnitude fainter than that of sources detectable in the 2 Ms *Chandra* survey.

Some of the energy absorbed by the circumnuclear obscuring material is reprocessed and reemitted at longer wavelengths, allowing for detection in the infrared (IR). NGC 1068 would have a flux density of $25\mu\text{Jy}$ at $24\mu\text{m}$ if placed at $z = 2$. While very faint, it would still be detectable in the ultradeep *Spitzer* MIPS (Rieke et al. 2004) image taken as part of the Great Observatories Origins Deep Survey (GOODS). The difficulty faced in mid-IR surveys is that the AGNs are vastly outnumbered by normal galaxies and some selection method is necessary to separate the two populations. Selection criteria which use *Spitzer*/IRAC mid-IR colors (Lacy et al. 2004; Stern et al. 2005) appear to be more prone to residual galaxy contamination at faint optical magnitudes or faint X-ray fluxes. Mid-IR spectral energy distribution (SED) techniques (Alonso-Herrero et al. 2006; Polletta et al. 2006; Donley et al. 2007) are more successful in picking out AGNs, but not necessarily the bulk of the CT AGN population.

Barger et al. (2007) used a highly spectroscopically complete deep VLA survey of the HDF-N region to examine whether radio sources can account for the missing 4 – 8 keV light. They found that the X-ray faint and X-ray undetected radio sources contributed only a few percent of the light at these energies and cannot account for the background light missing at these energies.

Optical narrow emission lines offer an intriguing alternative route to identifying high-redshift CT AGNs. As discussed in the previous section, since the narrow line region lies outside of the dusty circumnuclear material, the lines should not suffer from obscuration by that high column density material. Optical emission line ratio diagnostics separate AGNs from star-forming galaxies. A highly optically spectroscopically complete sample of X-ray selected AGNs can be used to establish a reliable diagnostic that pushes to higher redshifts than the classic BPT diagram described in the previous section (which is limited to $z < 0.5$). With a highly optically spectroscopically complete sample of galaxies with overlying X-ray data, we can use these diagnostic diagrams to pick out AGNs and then, comparing their predicted, intrinsic X-ray luminosities with their observed X-ray luminosity, identify candidate CT AGNs. The number density of CT AGNs determined using optical emission lines will provide an essential complementary view of the high-redshift CT

AGN population to the mid-IR and radio selection.

1.5 The OPTX Survey Fields

While other X-ray surveys have numerous redshifts, they are relatively incomplete in heterogeneous ways (see Table 2.1). This thesis focuses on three of the most uniformly observed and spectroscopically complete surveys to date. The deep pencil-beam *Chandra* Deep Field-North (CDF-N) survey is the most spectroscopically complete of all of the X-ray survey fields. Our group (Barger et al. 2002; Barger et al. 2003; Barger et al. 2005; Trouille et al. 2008) has spectroscopically observed 91% of the sources in this field.

Our two intermediate-depth, wide-field surveys – the *Chandra* Lockman Area Synoptic Survey (CLASXS) and the *Chandra* Lockman Area North Survey (CLANS) – provide an essential step between the ultradeep narrow *Chandra* surveys and the shallow wide-area surveys. They cover large cosmological volumes, detect rare, high-luminosity AGNs, and robustly probe AGN evolution between $z \sim 0$ and 1. Our group (Steffen et al. 2004; Trouille et al. 2008; Trouille et al. 2009) has spectroscopically observed 89% of the sources in the CLASXS field and 76% of the sources in the CLANS field.

Overall, we have spectroscopically observed 96% of the OPTX sources above $f_{2-8 \text{ keV}} = 10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2}$ and 78% below, predominantly using the DEIMOS and HYDRA spectrographs.

1.5.1 CDF-N

The 2 Ms CDF-N survey, which covers the well-known *Hubble* Deep Field-North and flanking fields, is currently the deepest ‘blank-field’ X-ray survey (although the *Chandra* Deep Field-South has recently been allotted a total of 4 Ms). The CDF-N has revealed 503 X-ray sources over a $\sim 460 \text{ arcmin}^2$ field-of-view to on-axis flux limits of $2.5 \times 10^{-17} \text{ erg cm}^{-2} \text{ s}^{-1}$ and $1.4 \times 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$ for the soft (0.5 – 2 keV) and hard (2 – 8 keV) X-ray bands, respectively (Alexander et al. 2003).

1.5.2 CLASXS & CLANS

The CLASXS and CLANS surveys sample large, contiguous solid angles (0.4 deg² and 0.6 deg², respectively) while remaining sensitive enough to measure 2 – 3 times fainter than the observed break in the 2 – 8 keV LogN-LogS distribution (Mushotzky et al. 2000). To minimize the effects of Galactic attenuation, CLASXS and CLANS are positioned in the Lockman Hole region of low HI galactic column density (Lockman et al. 1986).

The CLASXS survey consists of nine overlapping pointings with 40 ks exposures, except for the central pointing with 70 ks. This yields 525 X-ray sources with an on-axis 2 – 8 keV flux limit of $6 \times 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$. The CLANS survey consists of nine slightly overlapping 70 ks pointings. This yields 761 X-ray sources with an on-axis 2 – 8 keV flux limit of $3.5 \times 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$.

1.6 Thesis Goals and Outline

In this thesis I take advantage of the highly complete spectroscopic follow-up of the OPTX X-ray surveys to better understand the relation between the X-ray and optical spectral properties of AGNs. The

OPTX sample consists of 1789 sources in the CLANS, CLASXS, and CDF-N *Chandra* X-ray surveys covering 1.2 deg² of sky. The major goals of this thesis are to:

- Present the X-ray and optical photometric and spectroscopic data for the CLANS field and updated photometric and spectroscopic data for the CLASXS and CDF-N fields (Chapter 2).
- Compare the X-ray and optical spectral properties for a uniformly selected, highly spectroscopically complete sample of *Chandra* X-ray selected AGNs (Chapter 3).
- Determine the reliability of the [OIII] optical narrow emission line as an indicator of the ionizing flux from the central engine in AGNs (Chapter 4).
- Develop higher redshift optical narrow emission line ratio diagnostics that accurately identify X-ray selected AGNs (Chapter 5).
- Demonstrate that optical narrow emission line ratio diagnostics, when applied to galaxy samples with accompanying X-ray data, can be used to identify heavily X-ray absorbed, CT AGNs, providing an important alternative to mid-IR and radio searches (Chapter 5).

In Chapter 2 of this thesis I describe the creation of the X-ray catalog for the CLANS field and provide new (CLANS) and updated (CLASXS, CDF-N) redshift and optical/IR photometric catalogs for the three fields. In Chapter 3 I compare the optical spectral types with the X-ray spectral properties. In Chapter 4 I use the OPTX data to examine the reliability of optical emission line ratio diagnostics in identifying the distant X-ray AGN population and the reliability of the [OIII] emission line as an indicator of AGN activity. In Chapter 5 I use the OPTX sample to identify reliable high-redshift optical emission line ratio diagnostics. In applying these diagnostics to the highly optically spectroscopically complete sample of GOODS-N galaxies, I identify candidate CT AGNs and determine the number density for optically selected high-redshift CT AGNs. I give my conclusions, based on the results of my analyses, in Chapter 6.

This thesis encompasses all but one of the OPTX Project series articles (OPTX Project I, III, IV, and V). I have not included the analysis I carried out for the OPTX Project II article (Yenko et al. 2009). In this work, we combined the OPTX sample with the local *Swift* 9-month BAT sample to determine a local+distant hard X-ray luminosity function (HXLF) for the full sample and for the broad-line AGNs alone. We used a maximum likelihood method to find an analytic form for the HXLF, independent of the binning procedure, and presented the best-fit parameters for different luminosity and density dependent models.

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